

## Field-Induced Slow Magnetic Relaxation in Cobalt(II) Compounds with Pentagonal Bipyramid Geometry

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## Supporting Information

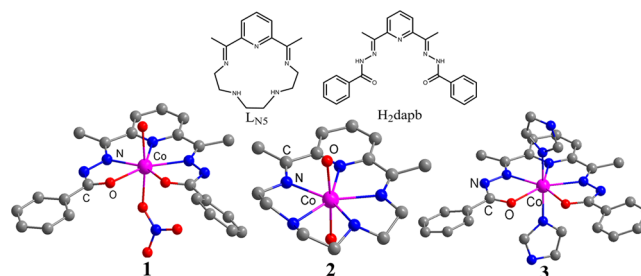
**ABSTRACT:** Field-induced slow magnetic relaxation was observed for air-stable mononuclear cobalt(II) compounds with pentagonal bipyramid geometry. These are the first examples of such behavior observed in the seven-coordinated mononuclear 3d metal compounds.

Single-molecule magnets (SMMs) continue to fascinate scientific communities thanks to their unique intrinsic magnetic properties and their possible applications in spintronic devices and quantum computing.<sup>1</sup> For the majority of SMMs containing 3d metal centers, the energy barrier  $U$  is proportional to the square of the ground state spin value  $S$ .<sup>2</sup> However, it was found to be extremely difficult to increase the barrier of SMMs despite the obvious success of synthesizing the metal clusters with very large  $S$  values. It is now a consensus that the high magnetic anisotropy is the most important to achieve better SMMs, although it is also the most difficult to design and control.

In this regard, the preparation of SMMs with a single metal center (single-ion magnet, SIM) has attracted intense interests. The most fascinating aspect of SIMs lies in the possible prediction and design of their magnetic anisotropy, and thus the SMM properties, based on the basic principles of the ligand field theory. In parallel with the fruitful SIMs based on lanthanides and actinides,<sup>3,4</sup> considerable efforts have also been devoted to the research of SIMs with 3d metal centers (3d-SIMs) since the first discovery of the Fe<sup>II</sup> SIM in 2010.<sup>5</sup> For all the reported 3d-SIMs, which are still quite limited, one common feature is the low coordination number of the metal centers, which affords a relatively weak ligand field to reduce the 3d orbital splitting energy, and to enhance the magnetic anisotropy.<sup>6</sup> With the coordination number ranging from 2 to 6, the most interesting 3d-SIMs contain the following metal centers: two-coordinated Fe<sup>I</sup> and Ni<sup>I</sup> ions;<sup>7</sup> two-, three-, four-, and six-coordinated Fe<sup>II</sup> ions;<sup>5,8</sup> five-coordinated Fe<sup>III</sup> ion;<sup>9</sup> three-, four-, five-, and six-coordinated Co<sup>II</sup> ions,<sup>10,11</sup> and six-coordinated Mn<sup>III</sup> ion.<sup>12</sup>

Meanwhile, the seven-coordinated metal centers with pentagonal bipyramid geometry (point group  $D_{5h}$ ) exhibit very large anisotropy, as established by both experimental and theoretical studies.<sup>10d,13</sup> Supported by the magnetization measurements and *ab initio* calculations, the zero-field splitting (ZFS) parameter  $D$  value of the pentagonal bipyramidal Co<sup>II</sup> center was found to be positive, which reflects the easy-plane magnetic anisotropy and seems to be unfavorable for a SIM at first glance. However, it has now been shown that slow magnetic relaxation is possible in the compounds of positive  $D$  values, as

observed recently in several Co<sup>II</sup> and Fe<sup>I/II</sup> compounds.<sup>7c,8a,10f–1,11</sup> For all of them, an external dc field was required to observe the slow magnetic relaxation. Interestingly, observation of the field-induced slow magnetic relaxation in a pentagonal bipyramid Co<sup>II</sup> compound [Co<sup>II</sup>(H<sub>2</sub>dapb)(H<sub>2</sub>O)(NO<sub>3</sub>)](NO<sub>3</sub>) (**1**) (H<sub>2</sub>dapb = 2,6-diacetylpyridine bis(benzoyl hydrazine), Figure 1) was actually suggested in the theoretical



**Figure 1.** Structures of  $L_{N5}$  and  $H_2dapb$  ligands and Co<sup>II</sup> **1–3**. Hydrogen atoms are omitted for clarity.

work by Ruiz et al., but no ac measurements has been reported for **1**.<sup>10d</sup> Very recently, Ruiz and Luis et al. proposed that the field-induced slow magnetic relaxation in the Kramers ions is a general consequence of the time-reversal symmetry that hinders the direct spin–phonon process.<sup>11</sup> Inspired by these studies, we reported here the dynamic magnetization study of compounds **1–3** ([Co<sup>II</sup> $L_{N5}$ (H<sub>2</sub>O)<sub>2</sub>]Cl<sub>2</sub>·4H<sub>2</sub>O (**2**) and [Co<sup>II</sup>(dapb)(im)<sub>2</sub>]·H<sub>2</sub>O (**3**),  $L_{N5}$  = 2,13-dimethyl-3,6,9,12-tetraaza-1(2,6)-pyridinacyclotridecaphane-2,12-diene, im = imidazole, Figure 1). Despite the different coordination environments of the Co<sup>II</sup> centers, field-induced slow magnetic relaxation was observed in all of them. This constitutes the very first observation of field-induced slow relaxation in compounds with pentagonal bipyramid geometry.

Air-stable compounds **1–3** can be synthesized conveniently (see details in the Supporting Information). The structure and dc magnetic property of **1** have been remeasured and are consistent with the reported results.<sup>13d</sup> A single-crystal X-ray diffraction study reveals that compound **2** crystallizes in the monoclinic space group  $P2_1/n$  and compound **3** in orthorhombic space group  $C222_1$  (Table S1). Their molecule structures are plotted in Figure 1. As shown, the Co<sup>II</sup> centers are in a pentagonal bipyramid environment with seven coordinated atoms, five of which are from equatorial ligands ( $N_5$  for **2** and  $N_3O_2$  for **1** and

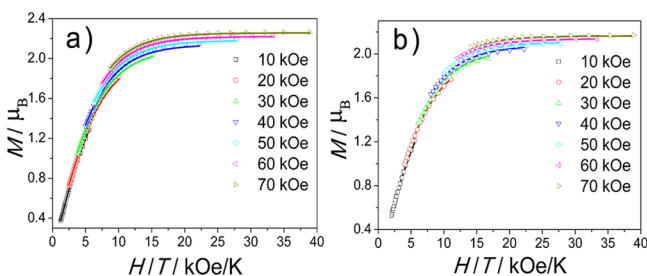
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3) and two of which are from axially coordinated water, nitrate ion, and imidazole molecules. The continuous shape measures (CShM's), used for calibrating the deviation of structures from a reference polygon, were calculated to be 0.417, 0.289, and 0.364 for **1**–**3**, which are close to zero of the ideal  $D_{5h}$  symmetry.<sup>14</sup>

Variable-temperature dc magnetic susceptibilities were measured on powder samples embedded in eicosane (Figure S1). The  $\chi_M T$  values at 300 K for **1**, **2**, and **3** are 2.63, 2.61, and 2.48  $\text{cm}^3 \text{mol}^{-1} \text{K}$ , respectively, which are larger than the spin-only value (1.875  $\text{cm}^3 \text{mol}^{-1} \text{K}$ ) for a high-spin  $\text{Co}^{\text{II}}$  ion. These values fall well in the range of 2.1–3.4  $\text{cm}^3 \text{mol}^{-1} \text{K}$  for the highly anisotropic  $\text{Co}^{\text{II}}$  ions with a considerable contribution from the orbital angular momentum.<sup>15</sup> Upon cooling from 300 K, the  $\chi_M T$  values of all these compounds decrease monotonously down to 1.8 K, reaching 1.56, 1.53, and 1.37  $\text{cm}^3 \text{mol}^{-1} \text{K}$  for **1**, **2**, and **3**, respectively. The decrease of the  $\chi_M T$  curves at low temperature is possibly due to the intrinsic magnetic anisotropy of the  $\text{Co}^{\text{II}}$  ions. The field dependence of the magnetization for **1**–**3** was carried out at 1.8 K ranging from 0 to 70 kOe (Figure S2). The magnetization at 70 kOe reaches 2.24, 2.26, and 2.16  $\mu_B$  for **1**, **2**, and **3**, respectively. The lack of saturation for the compounds also suggests the presence of appreciable magnetic anisotropy.

The reduced magnetization of **1**–**3** was measured at low temperatures with dc fields up to 70 kOe (Figure 2, Figure S3). A



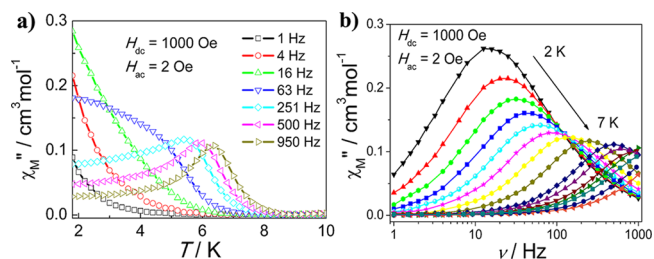
**Figure 2.** Reduced magnetization data for **2** (a) and **3** (b) collected under various applied dc fields. Solid lines represent the best fit using Anisofit 2.0.

spin Hamiltonian of eq 1 is utilized to describe the magnetic anisotropy qualitatively

$$H = D[S_z^2 - S(S+1)/3] + E(S_x^2 - S_y^2) + g\mu_B S \times B \quad (1)$$

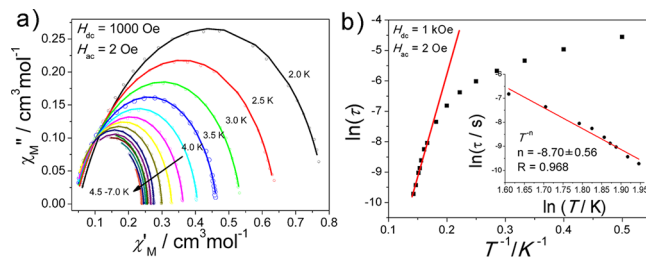
where  $\mu_B$  is the Bohr magneton, and  $D$ ,  $E$ ,  $S$ , and  $B$  represent single-ion axial and rhombic ZFS parameters, spin ( $S = 3/2$  for **1**–**3**), and magnetic field vectors, respectively. The best fits of the reduced magnetization data using Anisofit 2.0<sup>16</sup> give  $D = 32.4 \text{ cm}^{-1}$ ,  $E = 0$ , and  $g = 2.33$  for **1**;  $D = 24.6 \text{ cm}^{-1}$ ,  $E = -1.4 \times 10^{-2} \text{ cm}^{-1}$ , and  $g = 2.29$  for **2**;  $D = 24.8 \text{ cm}^{-1}$ ,  $E = 1.6 \times 10^{-3} \text{ cm}^{-1}$ , and  $g = 2.21$  for **3**. The results of  $D$ ,  $E$ , and  $g$  values of **1** agree well with the reported value of  $D = 31.0 \text{ cm}^{-1}$ ,  $E = 0$ , and  $g = 2.22$ .<sup>13b,d</sup> Using a negative initial value of  $D$ , another set of values of  $D \approx -14 \text{ cm}^{-1}$  and  $E \approx 11 \text{ cm}^{-1}$  for **1**–**3** were obtained with a poorer fit. These results conflict with the reported EPR data for **1**, indicating the correct choice of a positive sign of  $D$ .<sup>13d</sup> Furthermore, a positive  $D$  value was also reported in other similar  $\text{Co}^{2+}$  compounds with pentagonal bipyramid geometry<sup>13b,e</sup> and was demonstrated by theoretical calculations.<sup>10d,13d</sup> The positive sign of the  $D$  value stems from the interaction between the ground and excited electronic states coupled through spin–orbit coupling. However, we did not find any conclusive dependence of the  $D$  values on any parameters of the local geometries.

To probe the magnetic relaxation dynamics of all three compounds, temperature- and frequency-dependent ac susceptibilities were measured in the temperature range of 1.8–10 K. Under a zero dc field, no out-of-phase ac susceptibility ( $\chi''$ ) signal was observed for them (Figures S4, S7, S12). However, upon application of a 1000 Oe dc field, all compounds display temperature- and frequency-dependent ac signals typically observed for field-induced 3d-SIM species (Figure 3, Figures



**Figure 3.** Variable-temperature (a) and variable-frequency (b) out-of-phase ac magnetic susceptibility data for **1** in a 1000 Oe dc field.

S5–S6, S8–S9, S13–S14). In the temperature-dependent ac plots, the highest peaks of the  $\chi''$  at a frequency of 950 Hz are at 6.4, 5.8, and 6.2 K for **1**–**3**, respectively. Furthermore, the Cole–Cole plots below 7 K (Figure 4a, Figures S10, S15) were



**Figure 4.** (a) Cole–Cole plots of **1** under 1000 Oe dc field and (b) Arrhenius plot of  $\ln(\tau)$  vs  $T^{-1}$  for **1**. Solid line represents a fit to the high temperature region. Inset: Power law analysis in the form  $\ln(\tau)$  vs  $\ln(T)$ .

generated from the frequency-dependent ac susceptibility data. A generalized Debye model<sup>17</sup> was used to extract the values and distribution of the relaxation time ( $\tau$ ). The obtained  $\alpha$  values are 0.02–0.21, 0.01–0.26, and 0.03–0.14 for **1**–**3**, respectively. These values suggest the narrow distribution of the relaxation time. To estimate the effective energy barrier ( $U_{\text{eff}}$ ), Arrhenius plots [ $\ln(\tau)$  vs  $T^{-1}$ ] were constructed and plotted in Figure 4b and Figures S11 and S16. A linear fit to the high temperature data according to  $\tau = \tau_0 \exp(U_{\text{eff}}/k_B T)$  affords  $U_{\text{eff}} = 56.3 \text{ cm}^{-1}$  (81.2 K),  $20.7 \text{ cm}^{-1}$  (29.8 K), and  $62.3 \text{ cm}^{-1}$  (89.6 K) with  $\tau_0 = 6.0 \times 10^{-10} \text{ s}$ ,  $1.2 \times 10^{-6} \text{ s}$ , and  $8.7 \times 10^{-11} \text{ s}$  for **1**–**3**, respectively.

For a magnetic system with a positive  $D$  value, the mechanisms responsible for the slow magnetic relaxation are still a matter of debate. One of them arises from a transverse anisotropy barrier located in the easy plane, and the energy barrier is governed by a considerable  $E$  value.<sup>10g</sup> For our compounds, this mechanism does not seem applicable as the  $E$  values obtained experimentally (and as reported in other similar compounds with pentagonal bipyramid geometry) are very close to zero. The other mechanism was ascribed to a field-induced phonon bottleneck effect of the direct relaxation of the ground  $M_S = \pm 1/2$  levels. As a matter of fact, suggested by the very recent work of Ruiz and Luis,<sup>11</sup> direct phonon-induced processes are strongly suppressed

in a Kramers system with a considerable anisotropy irrespective of the sign of  $D$ . Thus, magnetic relaxation has to proceed either by the Orbach relaxation pathway through the excited  $M_S = \pm 3/2$  levels<sup>10f</sup> or by the optical acoustic Raman process involving a virtual state.<sup>10h,11</sup> For compound **2**, as the obtained energy barrier ( $20.7 \text{ cm}^{-1}$ ) is much lower than the energy gap between the  $M_S = \pm 1/2$  and  $M_S = \pm 3/2$  doublets, the Orbach pathway is not very likely. The relaxation times for **2** in the temperature range of 4–6 K can be fitted to a  $T^{-n}$  law with  $n = 4.9$  (Figure S11), which is close to the value of 4.5 for the Co–Y SIM reported by Colacio et al.,<sup>10h</sup> suggesting a dominant optical acoustic Raman process for the spin relaxation. However, for **1** and **3**, the obtained energy barriers are close to the 2D values. At the same time, the power law treatment of the data for **1** and **3** in the temperature range of 5–7 K gave  $n = 8.7$  (for **1**, Figure 4b) and 11.2 (for **3**, Figure S16), which is close to the expected  $n = 9$  for Raman relaxation in Kramers ions.<sup>18</sup> These results suggest that the Orbach and Raman processes seem possible for **1** and **3**.

In conclusion, our results demonstrate that slow magnetic relaxation can be observed under an applied dc field in the high-spin hepta-coordinate mononuclear Co<sup>II</sup> compounds with pentagonal bipyramid symmetry. This field-induced slow relaxation seems to be a quite common phenomenon for a Kramers compound with an easy-plane magnetic anisotropy. Studies on the dynamic magnetic properties of other metal centers, such as Fe<sup>II</sup>, Cr<sup>III</sup>, and Mn<sup>III</sup>, in the similar pentagonal bipyramid geometry are underway.

## ■ ASSOCIATED CONTENT

### ■ Supporting Information

X-ray crystallographic files (CIFs), experimental details, crystallographic data, and additional structural and magnetic figures. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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### Notes

The authors declare no competing financial interest.

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